## Other Optimization Techniques

## Conjugate Gradient

Similar to steepest descent, but slightly different way of choosing direction of next step:

$$
\begin{aligned}
& \vec{x}_{r+1}=\vec{x}_{r}+\lambda_{r} \vec{s}_{r} \\
& \vec{s}_{0}=-\vec{g}_{0} \\
& \vec{s}_{r+1}=-\vec{g}_{r+1}+\underbrace{\beta_{r+1} \vec{s}_{r}}_{\text {new term }}
\end{aligned}
$$

$\lambda_{r}$ is chosen to minimize $h\left(\vec{x}_{r+1}\right)$. This yields $\vec{g}_{r+1}^{\mathrm{t}} \vec{g}_{r}=0$
Here we allow a further step in the $\vec{s}_{r}$ direction. One choice
(Fletcher - Reeves) for $\beta_{r+1}$ is

$$
\beta_{r+1}=\frac{g_{r+1}^{2}}{g_{r}^{2}}
$$

## Newton-Raphson

Assume the function that we want to minimize is twice differentiable. Then, a Taylor expansion gives

$$
h(\vec{x}+\vec{\alpha}) \approx a+\vec{b}^{t} \vec{\alpha}+\frac{1}{2} \vec{\alpha}^{t} C \vec{\alpha}
$$

where

$$
a=h(\vec{x}), \quad \vec{b}=\vec{\nabla} h(\vec{x})=\vec{g}(\vec{x}), \quad C=\left(\frac{\partial^{2} h}{\partial x_{i} \partial x_{j}}\right)=H
$$

Now

$$
\vec{\nabla} h(\vec{x}+\vec{\alpha}) \approx \vec{b}+C \vec{\alpha}
$$

Because C is symmetric (check)

For an extremum, we have $\vec{b}+C \vec{\alpha}=0 \quad \vec{\alpha}=-C^{-1} \vec{b}$
or

$$
\vec{x}_{r+1}=\vec{x}_{r}-H\left(\vec{x}_{r}\right)^{-1} \vec{g}\left(\vec{x}_{r}\right)
$$

## Newton-Raphson

$$
\vec{x}_{r+1}=\vec{x}_{r}-H\left(\vec{x}_{r}\right)^{-1} \vec{g}\left(\vec{x}_{r}\right)
$$

i.e., the search direction is $\vec{s}=H\left(\vec{x}_{r}\right)^{-1} \vec{g}\left(\vec{x}_{r}\right) \quad$ and $\quad \lambda=1$

This converges quickly (if you start with a good guess), but the penalty is that the Hessian needs to be calculated (usually numerically)

Again, convergence is when $\vec{s}$ is sufficiently small
How would we calculate the Hessian numerically ? Use Lagrange polynomial in several dimensions and work it out

## Bounded Regions

The standard tool for minimization in particle physics is the MINUIT program (CERN library). It has also made its way well outside the particle physics community.

Author: Fred James

Here is how MINUIT handles bounded search regions - it transforms the parameter to be optimized as follows:

$$
\lambda^{\prime}=\arcsin \left(2 \frac{\lambda-a}{b-a}-1\right) \quad \lambda=a+\frac{b-a}{2}\left(\sin \lambda^{\prime}+1\right)
$$

$\lambda$ is the exernal (user) parameter
$\lambda^{\prime}$ is the internal parameter
MINUIT is available within PAW, ROOT ...

## MINUIT

MINUIT uses a (variable metric) conjugate gradient search algorithm (along with others). Basic idea:

- assume that the function to minimize can be approximated by a quadratic form near the minimum
- build up iteratively an approximation for the inverse of the Hessian matrix. Recall

$$
h(\vec{x}+\vec{\alpha}) \approx h(\vec{x})+\vec{\nabla} h(\vec{x}) \cdot \vec{\alpha}+\frac{1}{2} \vec{\alpha}^{t} H \vec{\alpha}
$$

the approximation for the Hessian is updated as follows:
$H_{i+1}=H_{i}+\frac{\left(\vec{x}_{i+1}-\vec{x}_{i}\right) \otimes\left(\vec{x}_{i+1}-\vec{x}_{i}\right)}{\left(\vec{x}_{i+1}-\vec{x}_{i}\right) \cdot\left(\vec{\nabla} h_{i+1}-\vec{\nabla} h_{i}\right)}-\frac{\left[H_{i} \cdot\left(\vec{\nabla} h_{i+1}-\vec{\nabla} h_{i}\right)\right] \otimes\left[H_{i} \cdot\left(\vec{\nabla} h_{i+1}-\vec{\nabla} h_{i}\right)\right]}{\left(\vec{\nabla} h_{i+1}-\vec{\nabla} h_{i}\right) \cdot H_{i} \cdot\left(\vec{\nabla} h_{i+1}-\vec{\nabla} h_{i}\right)}$
where the $\otimes$ symbol represents an outer product of two vectors (a matrix) $(\vec{a} \otimes \vec{b})_{i j}=a_{i} b_{j}$

## Fourier Transforms

Fourier transforms are very important

- as a way of summarizing the data with a few parameters
- because the transform of the data is itself very interesting (e.g., power spectrum, momentum $\Leftrightarrow$ coordinate space representation,...)

$$
\begin{array}{ll}
H(f)=\int_{-\infty}^{\infty} h(t) e^{2 \pi i f t} d t & H(f) \text { frequency domain representation } \\
h(t)=\int_{-\infty}^{\infty} H(f) e^{-2 \pi i t t} d t & h(t) \text { time domain representation }
\end{array}
$$

Warning: there is no unanimity on $2 \pi$ factors in front of the integral. Often the angular frequency is used $\omega=2 \pi f$

$$
H(\omega)=\int_{-\infty}^{\infty} h(t) e^{i \omega t} d t \quad h(t)=\frac{1}{2 \pi} \int_{-\infty}^{\infty} H(\omega) e^{-i \omega t} d \omega
$$

## Fourier Transform

Fourier Transform is a linear operation:

- transform of the sum of two functions is the sum of the transforms
- the transform of a constant times a function is constant times the transform

$$
\begin{array}{ll}
h(t) \text { real } & H(-f)=[H(f)]^{*} \\
h(t) \text { imaginary } & H(-f)=-[H(f)]^{*} \\
h(t) \text { even } & H(-f)=H(f) \\
h(t) \text { odd } & H(-f)=-H(f) \\
h(t) \text { real, even } & H(f) \text { real, even } \\
h(t) \text { real,odd } & H(f) \text { imaginary, odd } \\
h(t) \text { imaginary, even } & H(f) \text { imaginary, even } \\
h(t) \text { imaginary,odd } & H(f) \text { real, odd }
\end{array}
$$

## Fourier Transform

Further properties:

$$
\begin{aligned}
& h(a t) \Leftrightarrow \frac{1}{|a|} H\left(\frac{f}{a}\right) \\
& \frac{1}{|b|} h\left(\frac{t}{b}\right) \Leftrightarrow H(b f) \\
& h\left(t-t_{0}\right) \Leftrightarrow H(f) e^{2 \pi i f_{0}} \\
& h(t) e^{-2 \pi i f_{0} t} \Leftrightarrow H\left(f-f_{0}\right)
\end{aligned}
$$

We are typically interested in the Fourier analysis of a discretely sampled data set. Define the time step (taken to be constant here) as $\Delta$. The sampling rate (frequency) is $1 / \Delta$. Define the samples as

$$
h_{n}=h(n \Delta) \quad n=\cdots,-3,-2,-1,0,1,2,3, \cdots
$$

## Nyquist frequency

$$
f_{c} \equiv \frac{1}{2 \Delta} \quad \text { Nyquist frequency }
$$

This is the highest frequency which can be resolved with a sampling frequency $f=1 / \Delta$. If a continuous function $h(t)$ is limited in frequency components to frequencies less than $f_{c}$, then $h(t)$ is completely determined by its samples $h_{n}$. It can then be written as follows:

$$
h(t)=\Delta \sum_{n=-\infty}^{\infty} h_{n} \frac{\sin \left[2 \pi f_{c}(t-n \Delta)\right]}{\pi(t-n \Delta)}
$$

However, if there are frequency components which are higher than $f_{c}$, then they will be spuriously moved in the range $f<f_{c}$ (aliasing).

## Example

## Data

Fit
Aliasing



10 Hz sampling


5 Hz sampling

Conditions are: Sine wave with $\mathrm{f}=4 \mathrm{~Hz}$, phase offset $\varphi=0.1$ i.e., $h(t)=\sin (\varphi+2 \pi f t)=\sin (0.1+8 \pi t)$

## Example

$$
\begin{aligned}
& H(f)=\int_{-\infty}^{\infty} \sin (0.1+8 \pi t) e^{2 \pi i f t} d t \\
& =\int_{-\infty}^{\infty} \sin (0.1) \cos (8 \pi t) e^{2 \pi i f t} d t+\int_{-\infty}^{\infty} \cos (0.1) \sin (8 \pi t) e^{2 \pi i f t} d t \\
& =\sin (0.1) \int_{-\infty}^{\infty} \frac{e^{8 \pi i t}+e^{-8 \pi i t}}{2} e^{2 \pi i f t} d t+\cos (0.1) \int_{-\infty}^{\infty} \frac{e^{8 \pi i t}-e^{-8 \pi i t}}{2} e^{2 \pi i f t} d t
\end{aligned}
$$

## Recall the relation:

$\int_{-\infty}^{\infty} e^{2 \pi i f x} d f=\delta(x)$ where $\delta(x)$ is the Dirac Delta function
so we have

$$
\begin{aligned}
& H(f)=\sin (0.1) \int_{-\infty}^{\infty} \frac{e^{8 \pi i t}+e^{-8 \pi i t}}{2} e^{2 \pi i f t} d t+\cos (0.1) \int_{-\infty}^{\infty} \frac{e^{8 \pi i t}-e^{-8 \pi i t}}{2} e^{2 \pi i t t} d t \\
& =\frac{\sin (0.1)}{2} \int_{-\infty}^{\infty} e^{2 \pi i t(f+4)}+e^{2 \pi i t(f-4)} d t+\frac{\cos (0.1)}{2} \int_{-\infty}^{\infty} e^{2 \pi i t(f+4)}-e^{2 \pi i t(f-4)} d t \\
& =\frac{\sin (0.1)}{2}[\delta(f+4)+\delta(f-4)]+\frac{\cos (0.1)}{2}[\delta(f+4)-\delta(f-4)]
\end{aligned}
$$

## Discrete Fourier Transform

Suppose we have N consecutive sampled points

$$
h_{k} \equiv h\left(t_{k}\right), \quad t_{k} \equiv k \Delta, \quad k=0,1,2, \cdots, N-1
$$

We can extract the amplitude for N frequency components since we have N data points. Define the frequency components as

$$
f_{n} \equiv \frac{n}{N}\left(\frac{1}{\Delta}\right) \quad n=-\frac{N}{2}, \ldots, \frac{N}{2} \quad \text { (take } \mathrm{N} \text { even) }
$$

Note: there are $\mathrm{N}+1$ frequencies, but we will find that the two at the ends are equal, so only N independent. Negative frequencies allows us to include sine and cosine terms. So

$$
H\left(f_{n}\right)=\int_{-\infty}^{\infty} h(t) e^{2 \pi i f_{n} t} d t \approx \sum_{k=0}^{N-1} h_{k} e^{2 \pi i f_{n} t_{k}} \Delta=\Delta \sum_{k=0}^{N-1} h_{k} e^{2 \pi i f_{k} k \Delta}=\Delta \sum_{k=0}^{N-1} h_{k} e^{2 \pi i k n / N}
$$

$H_{n} \equiv \sum_{k=0}^{N-1} h_{k} e^{2 \pi i k n / N}$

## Discrete Fourier Transform

## Discrete Fourier Transform

The discrete fourier transform does not depend on any dimensional parameters.

Note $\quad H_{-n}=H_{N-n} \quad\left[e^{2 \pi i k(N-n) / N}=e^{2 \pi i k} e^{-2 \pi i k n / N}=e^{-2 \pi i k n / N}\right]$
In particular $\quad H_{-N / 2}=H_{N / 2}$

We can therefore rewrite the sum as follows

$$
H_{n} \equiv \sum_{k=0}^{N-1} h_{k} e^{2 \pi i k n / N} \quad n=0, \cdots, N-1
$$

Discrete inverse Fourier transform

$$
h_{k}=\frac{1}{N} \sum_{n=0}^{N-1} H_{n} e^{-2 \pi i k n / N} \quad k=0, \cdots, N-1
$$

## Discrete Fourier Transform

* Get the discrete Fourier components
* 

```
    Do \(n=0,63\)
    \(\mathrm{Hn}(\mathrm{n}, 1)=0 . \mathrm{D} 0\)
    \(\mathrm{Hn}(\mathrm{n}, 2)=0 . \mathrm{DO}\)
    Do \(k=0,63\)
        \(\mathrm{Hn}(\mathrm{n}, 1)=\mathrm{Hn}(\mathrm{n}, 1)\)
\& +amplitude( \(\mathrm{k}, 1)^{*}\) dcos(twopi*\({ }^{*}{ }^{*} \mathrm{n} / 64\).)
\& -amplitude(k,2)*dsin(twopi*k*n/64.)
        \(\mathrm{Hn}(\mathrm{n}, 2)=\mathrm{Hn}(\mathrm{n}, 2)\)
\& \(\quad+\) amplitude \((k, 1)^{*} d s i n\left(t w o p i *{ }^{*}{ }^{*} n / 64.\right)\)
\& \(\quad+\) amplitude(k,2)*dcos(twopi*k*n/64.)
    Enddo
    Write (11,*) N,Hn(N,1),Hn(N,2)
Enddo
```

amplitude( $k, 1$ ) real components
amplitude( $k, 2$ ) imaginary components

## Discrete Fourier Transform

Let's try it out on our sine wave data: Recall, signal $f=4 \mathrm{~Hz}$
Here's the result (64 points):




Large components are:
$f_{25}=\frac{25}{64 \cdot 0.1}=3.9, \quad f_{26}=\frac{26}{64 \cdot 0.1}=4.06$
$f_{38}=f_{64-26}=f_{-26} \quad f_{39}=f_{64-25}=f_{-25}$
cos and sin needed because of phase offset.

## Discrete Fourier Transform

## Here's the inverse transform

* Now we try the inverse transform

```
    Do n=0,63
    amplitude(n,1)=0.D0
    amplitude(n,2)=0.D0
    Do k=0,63
        amplitude(n,1)=amplitude(n,1)
    &
        +Hn(k,1)*dcos(twopi***n/64.)
                            +Hn(k,2)*dsin(twopi*k*n/64.)
        amplitude(n,2)=amplitude(n,2)
            -Hn(k,1)*dsin(twopi****/64.)
            +Hn(k,2)*dcos(twopi*k*n/64.)
        Enddo
    Write (12,*) N,amplitude(n,1)/64.,amplitude(n,2)/64.
    Enddo
```



## Fast Fourier Transform

The discrete Fourier transforms as we described it requires a sum over N terms for each of the N components. I.e., the number of operations scales as $\mathrm{N}^{2}$. A large part of the success of Fourier transforms for analysis of electronic signals, optical images, x-ray tomography,..., results from the fact that a numerical algorithm was found which requires of order $\mathrm{Nlog}_{2} \mathrm{~N}$ operations - the socalled Fast Fourier Transform (FFT). Here is how it works:

$$
\begin{aligned}
& F_{k}=\sum_{j=0}^{N-1} e^{2 \pi i j k / N} f_{j}=\sum_{j=0}^{N / 2-1} e^{2 \pi i k(2 j) / N} f_{2 j}+\sum_{j=0}^{N / 2-1} e^{2 \pi i k(2 j+1) / N} f_{2 j+1} \\
& =\sum_{j=0}^{N / 2-1} e^{2 \pi i k j /(N / 2)} f_{2 j}+W^{k} \sum_{j=0}^{N / 2-1} e^{2 \pi i k j /(N / 2)} f_{2 j+1} \\
& =F_{k}^{e}+W^{k} F_{k}^{o} \quad \text { where } W \equiv e^{2 \pi i / N}
\end{aligned}
$$

What is won? The sums in the individual terms in the last line only have $1 / 2$ as many terms, and the same factor appears

## Fast Fourier Transform

To see how this works in detail, we take an explicit example of having 8 data points (taking a power of 2 is important! If you don't have enough data, pad with zeroes).

$$
\begin{aligned}
& F_{k}=F_{k}^{e}+W^{k} F_{k}^{o} \\
& \text { where } W \equiv e^{2 \pi i / N}, F_{k}^{e} \equiv \sum_{j=0}^{3} W^{2 k j} f_{2 j}, F_{k}^{o} \equiv \sum_{j=0}^{3} W^{2 k j} f_{2 j+1}
\end{aligned}
$$

Now use a binary representation for the index $k=4 k_{2}+2 k_{1}+k_{0}$ where the $k_{i}$ 's are 0,1 . Then,
$W^{2 k j}=\left(e^{2 \pi i / 8}\right)^{2\left(4 k_{2}+2 k_{1}+k_{0}\right) j}=e^{2 \pi i\left(k_{2}+k_{1} / 2+k_{0} / 4\right)}=e^{2 \pi i\left(k_{1} / 2+k_{0} / 4\right)}=W^{2 j\left(2 k_{1}+k_{0}\right)}$
i.e., the $\mathrm{k}_{2}$ bit is irrelevant. So,
$F_{k}=F_{\left(k_{1}, k_{o}\right)}^{e}+W^{k} F_{\left(k_{1}, k_{o}\right)}^{o}$
$F_{\left(k_{1}, k_{0}\right)}^{e} \equiv \sum_{j=0}^{3} W^{2 j\left(2 k_{1}+k_{\mathrm{o}}\right)} f_{2 j} \quad F_{\left(k_{1}, k_{\mathrm{o}}\right)}^{o} \equiv \sum_{j=0}^{3} W^{2 j\left(2 k_{1}+k_{\mathrm{o}}\right)} f_{2 j+1}$

## Fast Fourier Transform

Let's try again:
$F_{k}^{e}=\sum_{j=0}^{3} W^{2 j\left(2 k_{1}+k_{0}\right)} f_{2 j}=\sum_{j=0}^{1} W^{2(2 j)\left(2 k_{1}+k_{0}\right)} f_{2(2 j)}+\sum_{j=0}^{1} W^{2(2 j+1)\left(2 k_{1}+k_{0}\right)} f_{2(2 j+1)}$
$=\sum_{j=0}^{1} W^{2(2 j)\left(2 k_{1}+k_{0}\right)} f_{2(2 j)}+W^{2\left(2 k_{1}+k_{0}\right)} \sum_{j=0}^{1} W^{2(2 j)\left(2 k_{1}+k_{\mathrm{o}}\right)} f_{2(2 j+1)}$
$F_{k}^{e}=F_{k}^{e e}+W^{2\left(2 k_{1}+k_{0}\right)} F_{k}^{e o}$

$$
W^{4 j\left(2 k_{1}+k_{0}\right)}=\left(e^{2 \pi i / 8}\right)^{\left(8 k_{1}+4 k_{0}\right) j}=W^{2 j k_{0}}
$$

and $\quad F_{k}^{o}=F_{k}^{o e}+W^{2\left(2 k_{1}+k_{0}\right)} F_{k}^{o o}$
Can perform one more step:

$$
F_{k}^{e e}=F^{e e e}+W^{4 k_{0}} F^{e e o} \quad F^{e e e}=f_{0} \quad F^{e e o}=f_{4}
$$

The sums have disappeared!

## Fast Fourier Transform

The final pieces are:

$$
\begin{aligned}
& F_{k}^{e e}=F^{e e e}+W^{4 k_{0}} F^{e e o}=f_{0}+W^{4 k_{o}} f_{4} \\
& F_{k}^{e o o}=F^{e o e}+W^{4 k_{o}} F^{e o o}=f_{2}+W^{4 k_{0}} f_{6} \\
& F_{k}^{o e}=F^{o e e}+W^{4 k_{o}} F^{o e o}=f_{1}+W^{4 k_{o}} f_{5} \\
& F_{k}^{o o}=F^{o o e}+W^{4 k_{o}} F^{o o o}=f_{3}+W^{4 k_{o}} f_{7}
\end{aligned}
$$

Note $\mathrm{k}_{0}=0,1$
So, need 8 multiplications and 8 additions for this step

Then,

$$
\begin{aligned}
& F_{k}^{e}=F^{e e}+W^{2\left(2 k_{1}+k_{o}\right)} F^{e o} \\
& F_{k}^{o}=F^{o e}+W^{2\left(2 k_{1}+k_{0}\right)} F^{o o}
\end{aligned}
$$

Here $\mathrm{k}_{0}=0,1 \quad \mathrm{k}_{1}=0,1$ So, again 8 multiplications and 8 additions for this step

Finally $\quad F_{k}=F^{e}+W^{k} F^{o}$ again 8 multiplications and 8 additions for this step

## Fast Fourier Transform

So we need 2N operations per level, and there are $\log _{2} \mathrm{~N}$ levels. The scaling of the computational time is therefore $\mathrm{Nlog}_{2} \mathrm{~N}$ rather than $\mathrm{N}^{2}$.
E.g., $\mathrm{N}=1000 \mathrm{Nlog}_{2} \mathrm{~N} \sim 1000 * 10=10^{4} \quad \mathrm{~N}^{2}=10^{6}$

How to implement in practice. Note that the trick is to find out which value of $n$ corresponds to which pattern of e,o in

$$
F^{\text {eoeoоеое... }}=f_{n}
$$

Answer: reverse pattern of $e, o$. Assign $e=0,0=1$, and the binary value gives $n$.

$$
\text { eеe } \rightarrow \text { eеe } \rightarrow 000 \rightarrow 0
$$

$$
\text { oeo } \rightarrow \text { eoe } \rightarrow 010 \rightarrow 2
$$

## Some examples

. Agric. Food Chem., 52 (20), 6055 -6060, 2004. 10.1021/jf049240e S0021-8561(04)09240-4
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Discrimination of Olives According to Fruit Quality Using Fourier Transform Raman Spectroscopy and Pattern Recognition Techniques

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Multiplication of large integers
The fastest known algorithms for the multiplication of large integers or polynomials are based on the discrete Fourier transform: the sequences of digits or coefficients are interpreted as vectors whose convolution needs to be computed; in order to do this, they are first Fourier-transformed, then multiplied component-wise, then transformed back.

## Power Spectrum

The autocorrelation of a function is

$$
\operatorname{Corr}[y](\tau)=\int_{-\infty}^{\infty} y(t)^{*} y(t+\tau) d t
$$

and the power spectrum is defined as the Fourier transform of the autocorrelation

$$
P S[y](f)=\int_{-\infty}^{\infty} \operatorname{Corr}[y](\tau) e^{2 \pi i f \tau} d \tau
$$

For a periodic function, the correlation is often defined as the expectation value. There is no convention on the normalization, so be careful about the values. Best to see which frequencies dominate a given spectrum. Here is a practical approach for a discretely sampled function:

$$
C_{k}=\sum_{j=0}^{N-1} c_{j} e^{2 \pi i j k / N} \quad k=0,1, \ldots, N-1
$$

## Power Spectrum

$$
\begin{aligned}
& P(0)=P\left(f_{0}\right)=\frac{1}{N^{2}}\left|C_{0}\right|^{2} \\
& P\left(f_{k}\right)=\frac{1}{N^{2}}\left[\left|C_{k}\right|^{2}+\left|C_{N-k}\right|^{2}\right] \quad k=1,2, \ldots,\left(\frac{N}{2}-1\right) \\
& P\left(f_{c}\right)=P\left(f_{N / 2}\right)=\frac{1}{N^{2}}\left|C_{N / 2}\right|^{2}
\end{aligned}
$$

Let's try it out on our example:
The 4 Hz frequency is picked out.


## Noise

## We now add some noise to our spectrum (Gaussian smearing

 with $\sigma=0.5$ ) and see what happens:

Data sample


Winter Semester 2006/7



Data sample


Computational Physics Í

Inverse Fourier Transform


Inverse Fourier Transform


Lecture 1225

## Exercizes

1. Solve the $\chi^{2}$ minimization problem from last lecture with MINUIT.
2. Generate 64 data points using

$$
f(t)=\cos \left(\pi / 4+2 \pi f_{1} t\right)+\cos \left(2 \pi f_{2} t\right) \quad f_{1}=0.5, f_{2}=1 \quad \Delta=0.2
$$

and fit with a discrete Fourier transform. Extract the power spectrum.

