

# Experimental Review of Photon Structure Function Data



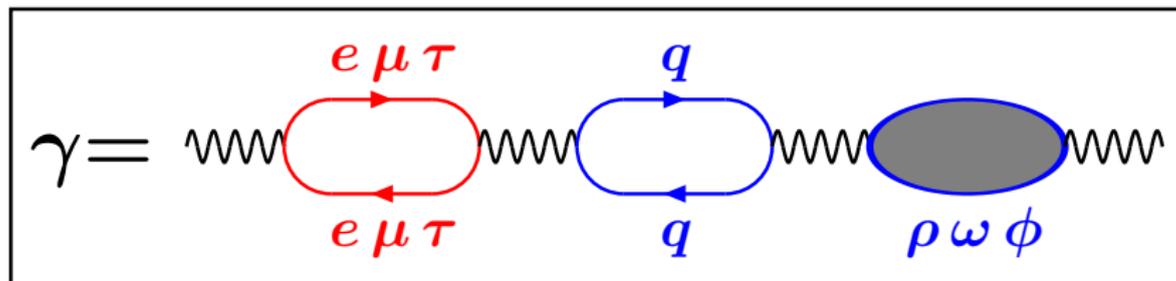
Hamburg, May 12, 2009

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## Why do we talk about photon structure?

- The structure of the photon is a purely quantum mechanical effect.



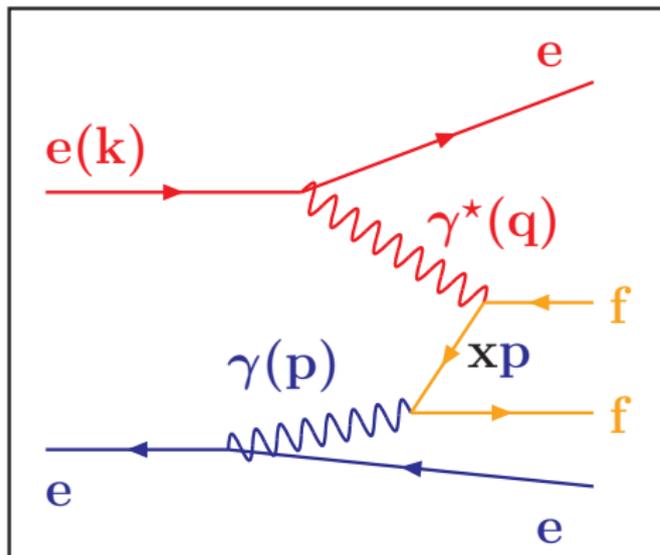
- Due to the Heisenberg uncertainty principle the photon may, for a short amount of time, fluctuate into a **leptonic** or **hadronic** state (with the same quantum numbers as the photon).
- The typical life time  $\Delta t = 1/\Delta E$  of these states **increases with the photon energy** and **decreases with the photon virtuality**.



$$\Delta p \cdot \Delta q \geq \frac{1}{2} \hbar$$

**The photon structure is enriched for quasi-real, high energetic photons.**

## How do we measure photon structure functions?



### Deep-inelastic Electron-Photon Scattering

$Q^2 = -q^2 \gg 0 \Rightarrow$  this electron is visible within the detector.

$$x = \frac{Q^2}{Q^2 + W^2}, \quad y = \frac{pq}{pk}$$

$P^2 = -p^2 \approx 0 \Rightarrow$  this electron stays within the beam pipe.

– The differential cross-section:

$$\frac{d^2\sigma}{dx dQ^2} \approx k(x, y, Q^2) \cdot F_2^\gamma(x, Q^2, P^2)$$

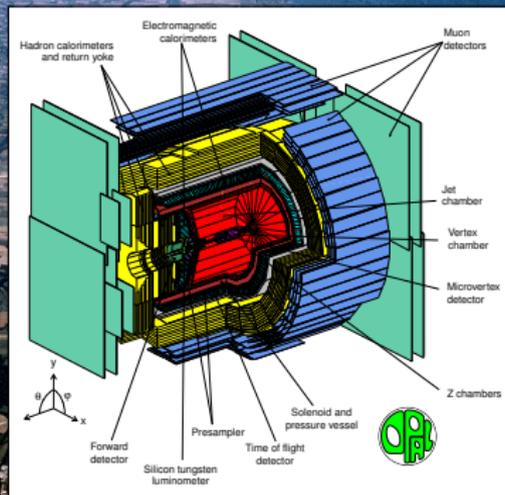
The structure function  $F_2^\gamma$  parametrises the internal structure of the photon.

## The 'history' of photon structure function measurements

Date	Event
1973	Investigation of two-photon processes in QPM by Walsh and Zerwas
1977	The LO asymptotic behavior of $F_2^\gamma \propto 1/\alpha_s$ was discovered by Witten
1979	Calculation of NLO corrections by Bardeen and Buras
1981	The first measurement of $F_2^\gamma$ by PLUTO
1986	The first extraction of $\Lambda$ from $F_2^\gamma$ data
1990	Start of $F_2^\gamma$ measurements at TRISTAN
1994	Start of $F_2^\gamma$ measurements at LEP
2002	NLO extraction of $\alpha_s$ based on a large set of data by Albino et. al
2005	The final LEP2 results are being published
2011	First measurement of $F_2^\gamma$ by Belle and Babar?
2018	First measurement of $F_2^\gamma$ at a future Linear Collider?

**A long tradition: Unfortunately the last two dates have to be changed from time to time.**

# The Large Electron Positron Collider (1989 - 2000): $E_{cm} = 90 - 209 \text{ GeV}$

 $e^-$  $e^+$ 

**Data statistics:**

160/pb at  $E_{cm} \approx 90 \text{ GeV}$

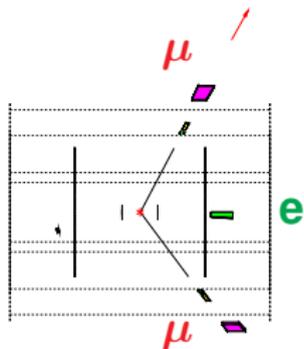
700/pb at  $E_{cm} > 160 \text{ GeV}$

## Two typical events

```
Run.event: 5198-223277 Date: 940625 Time: 211945 Ctx(Nr: 2 Surp: 7.3) Ecal(Nr: 3 SurEn: 1.4) Hcal(Nr: 4 SurEn: 3.3)
Ebeam: 45.62 Evls: 10.5 Emiss: 80.7 Vtx: (-0.62, 0.34, 0.47) Muon(Nr: 2) Sec Vtx(Nr: 0) Fast(Nr: 0 SurEn: 0.0)
Bv:4:028 Bunchlet: 11 Thrust:0.849 Aplan:0.973 Q:lat:0.4978 Spherr:0.4109
```

Event type bits

```
4 Low mult presel
12 Tagged two phot
22 S phot muon veto
30 "Physics" selection
1 ZD type physics
```



$e\gamma \rightarrow e\mu^+\mu^-$

Centre of screen is ( 0.0000, 0.0000, 0.0000)

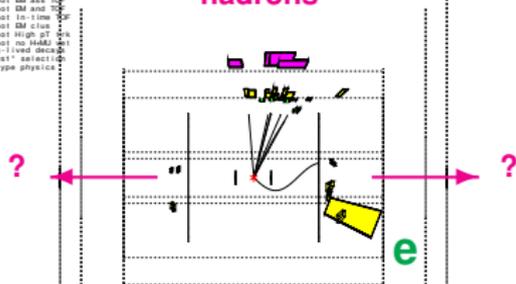
200 cm. 510 20 55 GeV

```
Run.event: 6422-47894 Date: 950817 Time: 155340 Ctx(Nr: 8 Surp: 12.4) Ecal(Nr: 19 SurEn: 46.8) Hcal(Nr: 6 SurEn: 3.4)
Ebeam: 45.64 Evls: 58.0 Emiss: 33.3 Vtx: (-0.85, 0.11, 1.11) Muon(Nr: 0) Sec Vtx(Nr: 0) Fast(Nr: 0 SurEn: 0.0)
Bv:4:028 Bunchlet: 213 Thrust:0.7845 Aplan:0.8856 Q:lat:0.4769 Spherr:0.8370
```

Event type bits

```
4 Low mult presel
8 Single phot presel
12 Tagged two phot
13 High pT high mult
24 S phot EM ass tag
25 S phot EM ass tag
26 S phot in-time tag
27 S phot EM cut
28 S phot High pT cut
29 S phot no HADM cut
31 long-lived deca
32 "Physics" selection
1 ZD type physics
```

hadrons



$e\gamma \rightarrow e q\bar{q}$

Centre of screen is ( 0.0000, 0.0000, 0.0000)

200 cm. 510 20 55 GeV



- The **scattered electron** and the **two muons** are clearly visible.
- Clean topology and good mass resolution.

- The **scattered electron** is clearly visible.
- The **hadronic final state** may partly disappear along the beam axis.

The hadronic final state is much harder to measure.

## The unfolding of $F_2^\gamma$ from the data

### The task

- Deduce the underlying function  $f(x)$  from the measured distribution:

$$g^{\text{det}}(x_{\text{vis}}, \text{Da}) = \int A(x_{\text{vis}}, x) f(x) dx + B(x_{\text{vis}})$$

### The solution

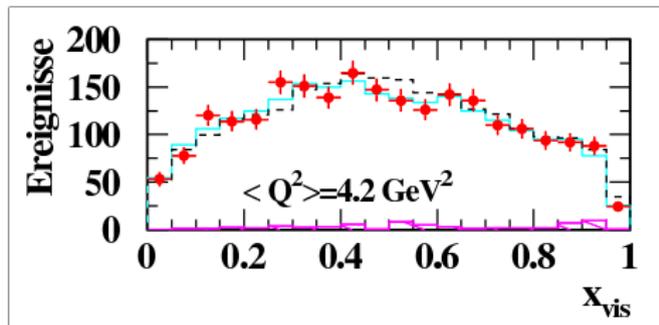
- Monte Carlo (MC) simulation of many events and unfolding of the distribution by:

- 1) Simulation of many MC signal events based on some  $\tilde{f}(x) \Rightarrow A(x_{\text{vis}}, x)$ .
- 2) Simulation of many MC background events  $\Rightarrow B(x_{\text{vis}})$ .
- 3) Solve the integral  $\rightarrow$  matrix equation numerically (with regularisation), i.e. fit the  $g^{\text{det}}(x_{\text{vis}}, \text{MC})$  to the data distribution  $g^{\text{det}}(x_{\text{vis}}, \text{Da})$  by variation of  $f(x) = \tilde{f}(x) \cdot c(x)$ .

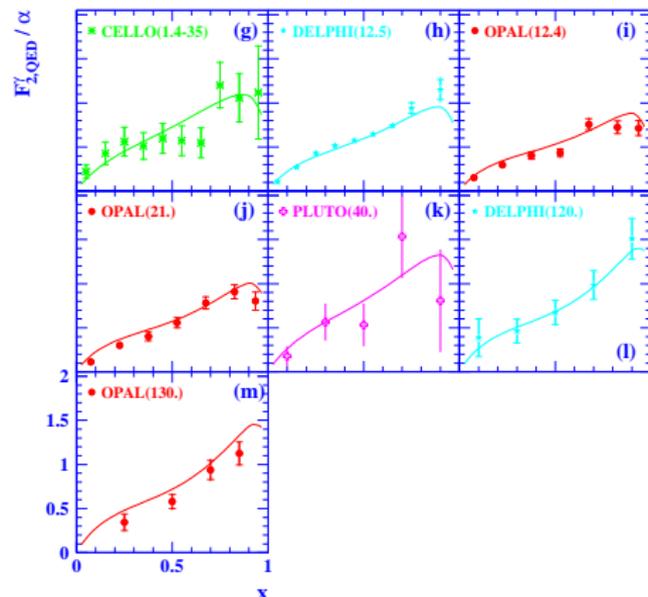
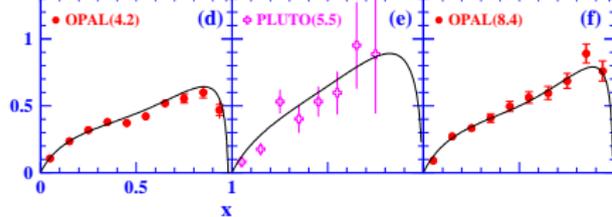
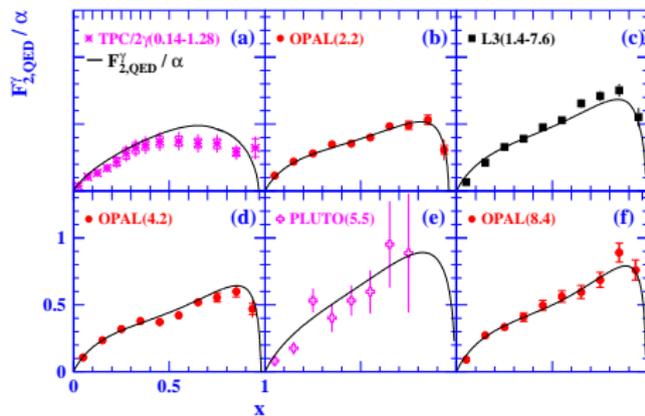
### The result

- After the fit, the distribution  $g^{\text{det}}(x_{\text{vis}}, \text{MC})$  and  $g^{\text{det}}(x_{\text{vis}}, \text{Da})$  are identical within errors, this means the structure function is:

$$F_2^\gamma(x, \text{Da}) = c(x) \cdot F_2^\gamma(x, \text{MC})$$



## The world data on $F_{2,QED}^\gamma$

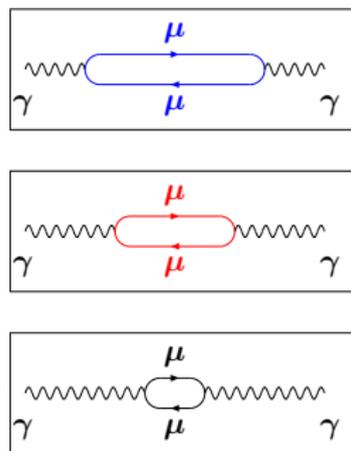
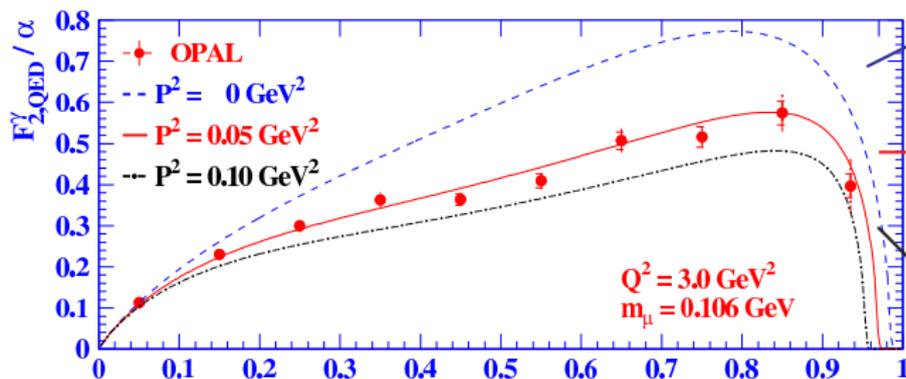


- The data covers the virtuality range  $2 < \langle Q^2 \rangle < 130 \text{ GeV}^2$ .
- The precision is a few per cent.
- There is even more to come, see talk by K. Dehmelt.

The data on  $F_{2,QED}^\gamma$  span two orders of magnitude in  $Q^2$  and are very precise.

## The structure of virtual photons is suppressed

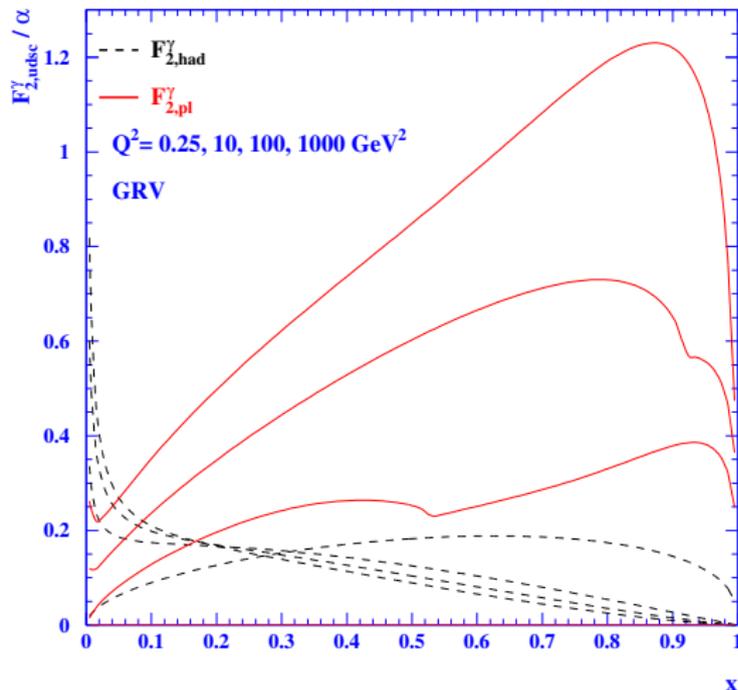
- The measurement of the reaction  $e \gamma \rightarrow e \mu^+ \mu^-$  gives:



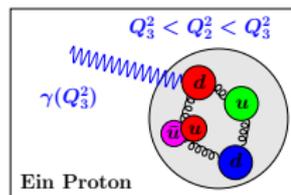
- For  $P^2 = 0$  the photon is real and  $F_{2,QED}^\gamma$  is maximal.
- The Standard Model prediction is  $P^2 = 0.05 \text{ GeV}^2$ .
- For  $P^2 \gg 0$  the photon is highly virtual and  $F_{2,QED}^\gamma$  is reduced.

The suppression of the photon structure for virtual photons is clearly seen in the data.

## The hadronic structure of the photon

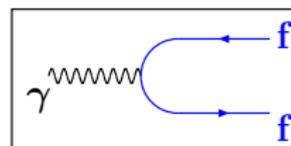


– The proton is a hadron.



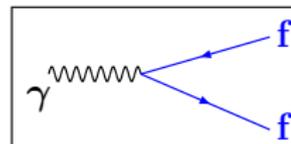
the structure depends on  $Q^2$ .

– The photon has a **hadron-like**



$\equiv$  proton

and a **point-like**

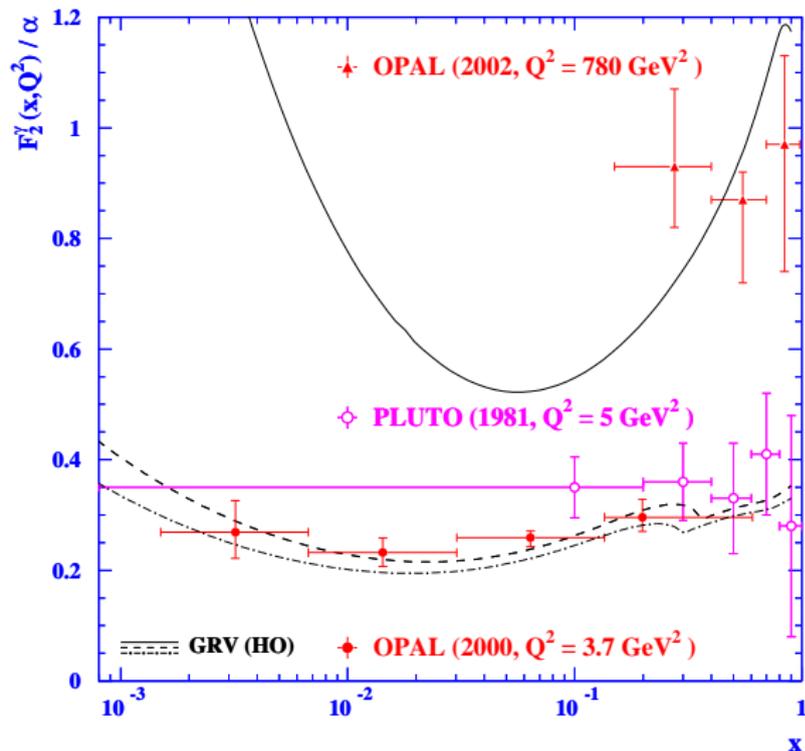


$\neq$  proton

component.

Due to the point-like component  $F_2^\gamma$  rises with  $Q^2$  for all values of  $x$ .

## What a difference 24 years make



### Kinematical range covered

- About a factor 100 decrease in  $x$
- About a factor 100 increase in  $Q^2$

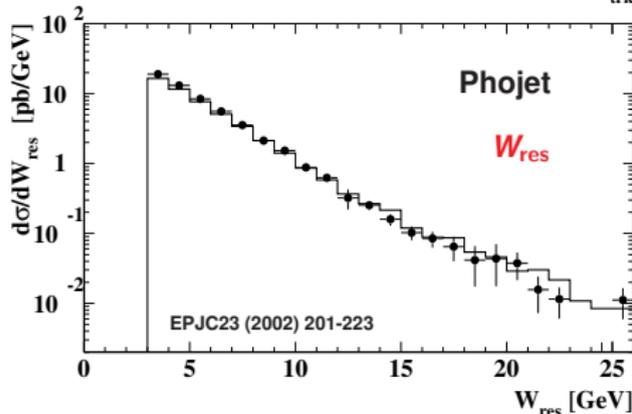
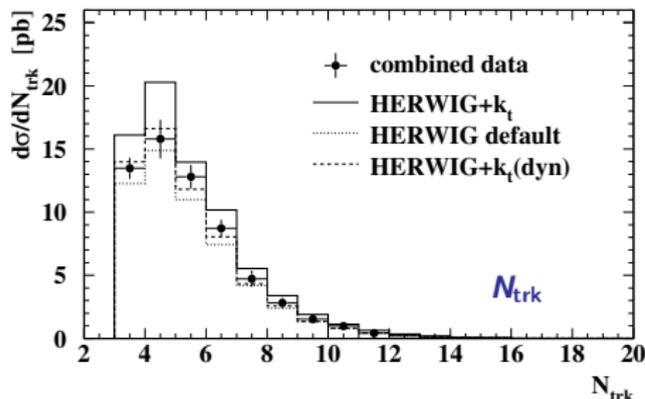
### Analysis methods

- Multipurpose MC models
- Radiative corrections
- Sophisticated unfolding methods
- LEP combined effort
- About 50 measurements



**Significantly smaller errors !?**

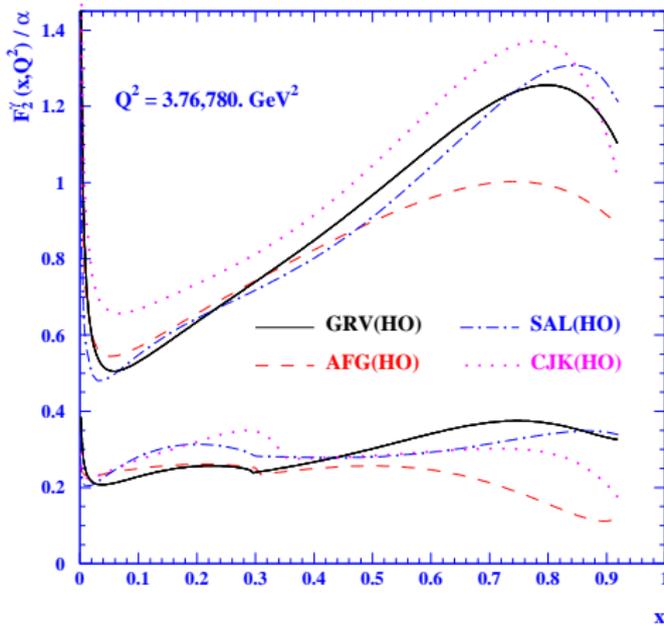
## The main problem when extracting $F_2^\gamma$



- The quality of the description of the hadronic final state by the Monte Carlo models is far from perfect.
- However, the LEP data show consistent deviations for several observables.
- The combination of the LEP data by the LEP 2 $\gamma$ -WG results in smaller uncertainties for observables like:
  - 1) the charged multiplicity  $N_{\text{trk}}$
  - 2) the observed invariant mass  $W_{\text{res}}$  in a restricted acceptance region
- This helps for adjustments of model parameters in collaboration with the authors of the Monte Carlo programs.

**Still, for large parts of the phase space this is the main systematic uncertainty.**

## Some recent higher order $F_2^\gamma$ parametrisations



- All groups have problems fitting the DELPHI prel. data  $\Rightarrow$  exclusion or error inflation.

A number of new parametrisations with different theoretical assumptions are available.

### The CJK(HO) parametrisation

- All  $F_2^\gamma$  data incl. TPC/2 $\gamma$  and DELPHI prel. (LEP1).
- $Q_0^2 = 0.765 \text{ GeV}^2$ ,  $\Lambda_4^{\overline{\text{MS}}} = 280 \text{ MeV}$ .

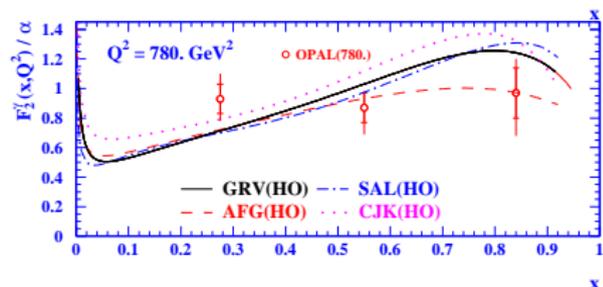
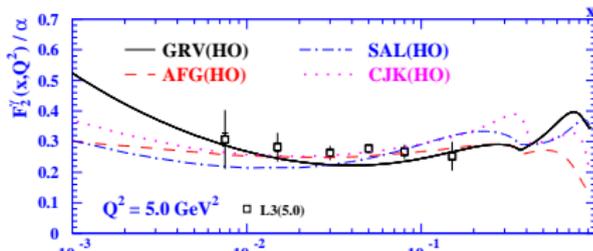
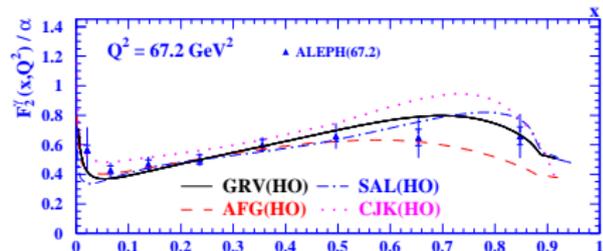
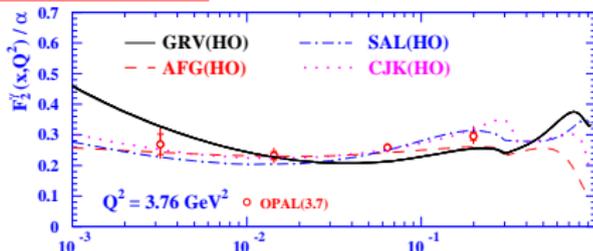
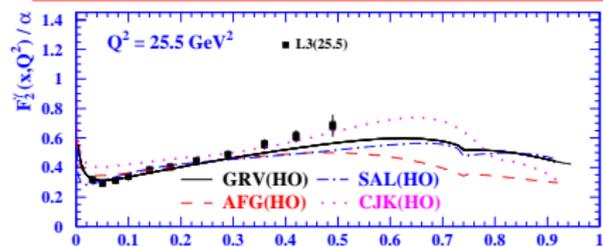
### The AFG(HO) parametrisation

- LEP1 data at medium  $Q^2$ , incl. DELPHI prel.
- Massless  $Q = c, b$ , but  $m_Q^2/Q^2$  corrections.
- $Q_0^2 = 0.7 \text{ GeV}^2$ ,  $\Lambda_4^{\overline{\text{MS}}} = 300 \text{ MeV}$ .

### The SAL(HO) parametrisation

- All  $F_2^\gamma$  data besides TPC/2 $\gamma$  + ZEUS  $F_2^p$  at  $x < 0.01$  + ZEUS dijet data.
- Gribov facto.:  $F_2^\gamma = \frac{\sigma_{\gamma p(W)}}{\sigma_{pp(W)}} \cdot F_2^p \approx 0.43 \cdot F_2^p$
- $Q_0^2 = 2.0 \text{ GeV}^2$ ,  $\Lambda_4^{\overline{\text{MS}}} = 330 \text{ MeV}$ .

## Recent parametrisations compared to $F_2^\gamma$ data

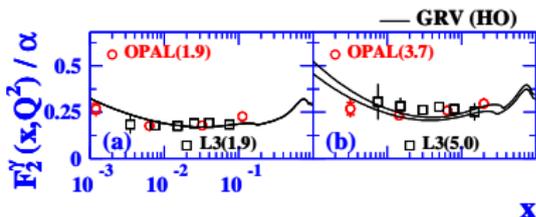


### A few observations

- The lowest  $Q^2$  data is best described by AFG.
- CJK seems too high at low- $x$  wrt. L3 data.
- At the highest  $Q^2$  AFG fits the data best.

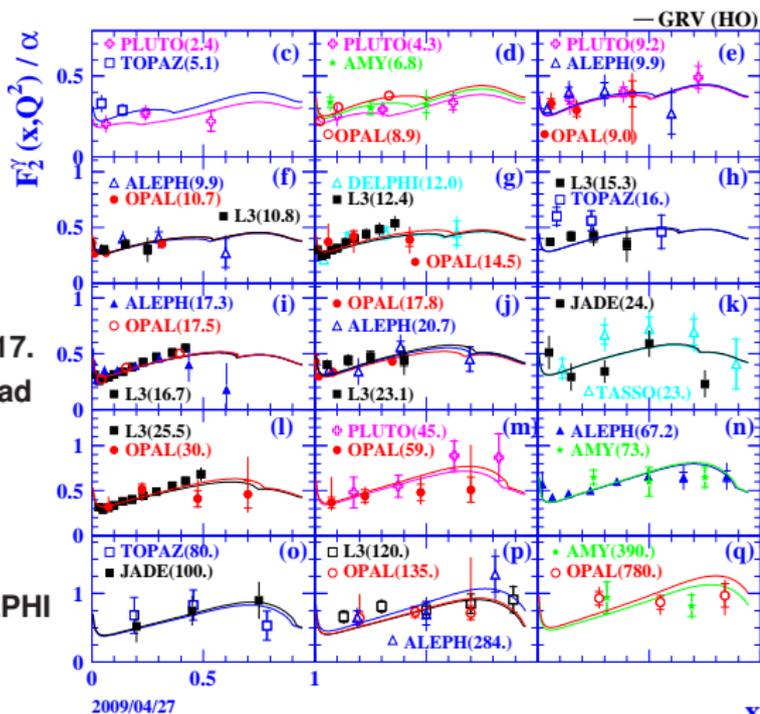
A complete survey would be needed to find the set that is preferred by the data.

## The world data on $F_{2,had}^\gamma$



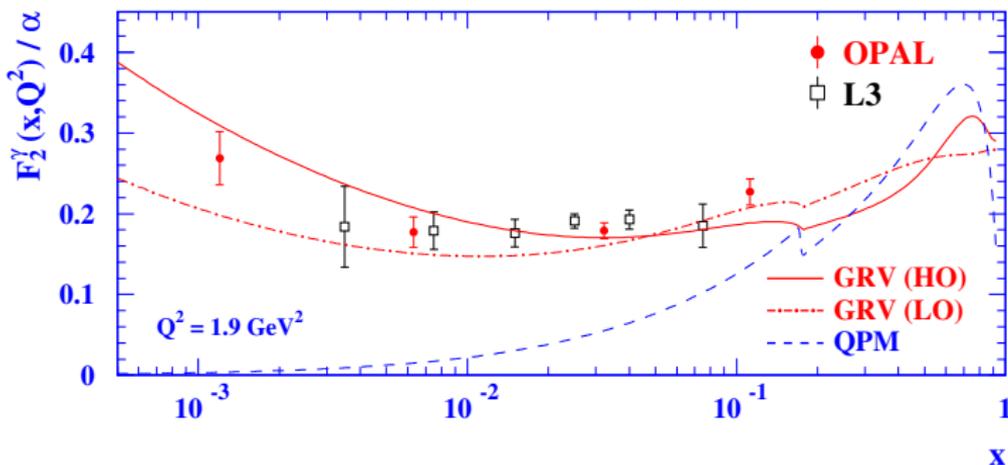
### Some details on the selection of data

- Start from Phys.Rep. 332 (2000) 165-317.
- Drop the TPC/ $2\gamma$  results due to very bad  $\chi^2$  wrt. various  $F_2^\gamma$  parametrisations (see Tables 4+5) in Phys.Rep. 332.
- Add newly published results from ALEPH and L3.
- Drop the preliminary results from DELPHI which did not get published by now.



Many measurements. The precision is dominated by the results from LEP.

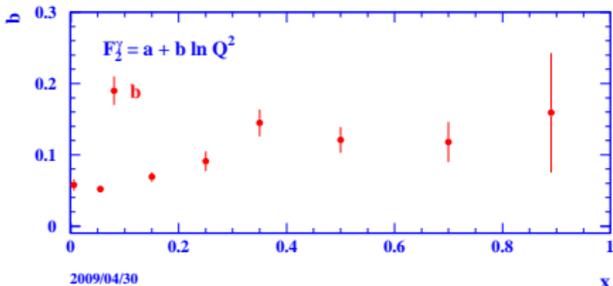
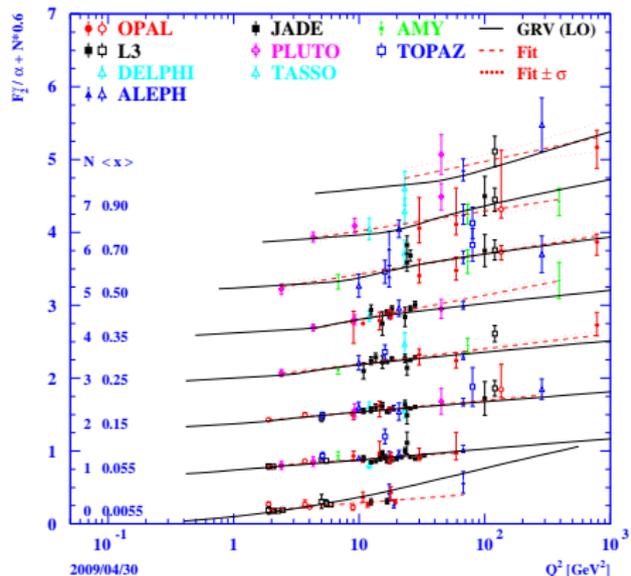
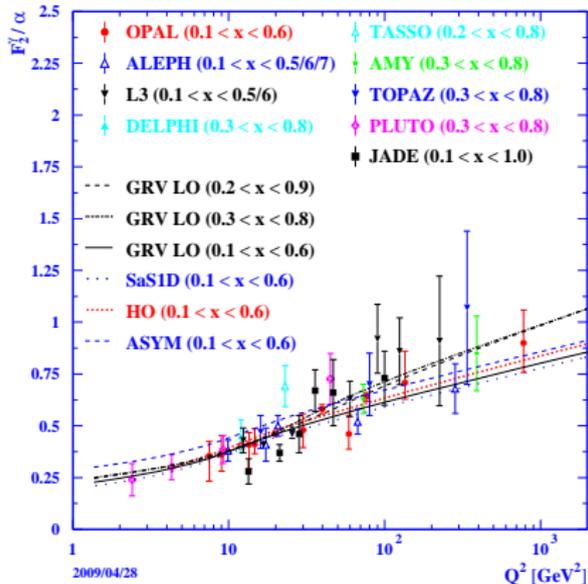
## The measurements of $F_2^\gamma$ at low $x$ and $Q^2$



- The LEP data are consistent and determine  $F_2^\gamma$  to 5-20% precision.
- The expected rise of  $F_2^\gamma$  is still very moderate.
- The QPM prediction is much too low compared to the data.
- QCD expectations, e.g. the GRV parametrisation are able to account for the data.

Unfortunately, the kinematical region is too small to test the low- $x$  rise.

## The $Q^2$ evolution of $F_2^\gamma$



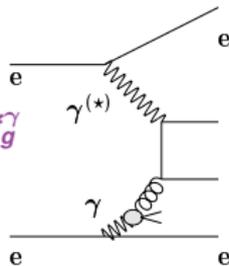
— This comparison has to be made by carefully selecting the  $x$  range.

$F_2^\gamma$  rises with  $Q^2$  for all ranges of  $x$ .

## The $F_{2,c}^\gamma$ measurement

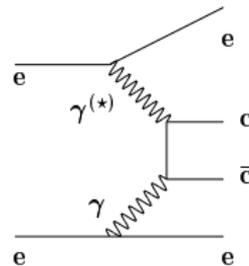
hadron-like:

- depends on  $f_g^\gamma$
- low- $x$

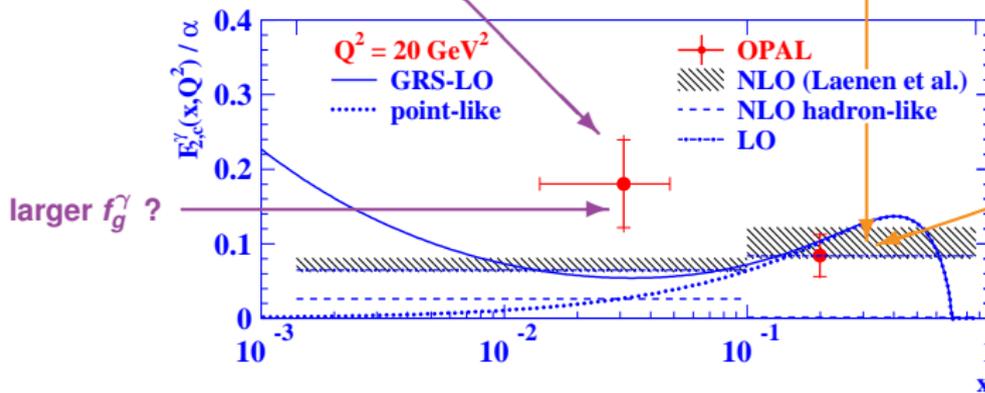
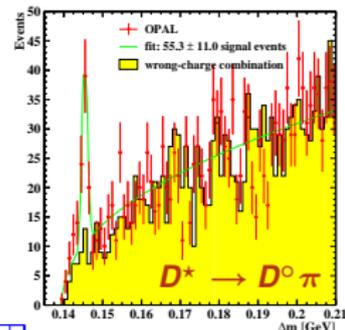


point-like:

- pQCD
- high- $x$



## Charm tagging

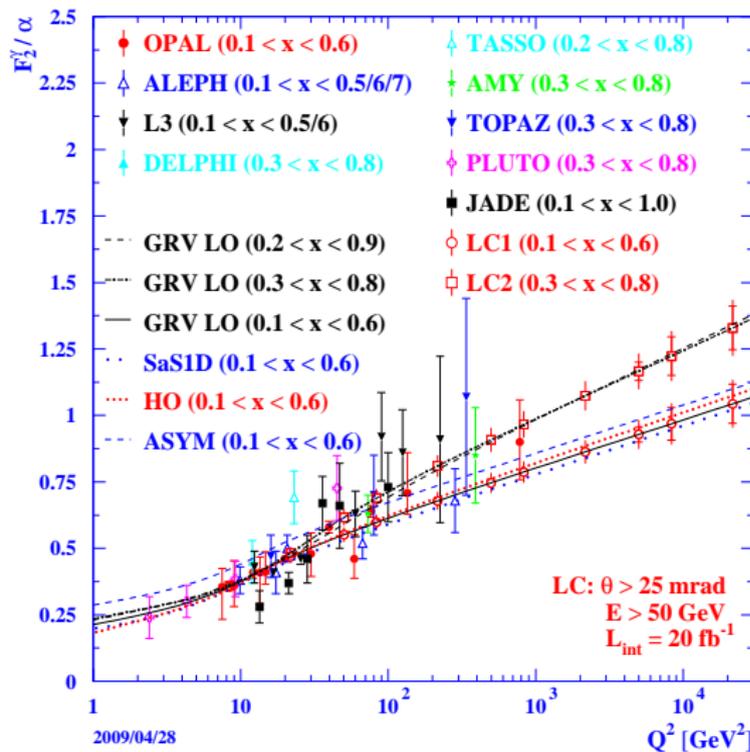


NLO =  $f(\alpha_s, m_c)$   
perfectly fits

PLB539 (2002) 13-24

**This result is not yet conclusive. It needs ALEPH, DELPHI and L3 to confirm .**

## The future of $F_2^\gamma$ measurements



### The assumptions for the ILC Data

–  $e^+e^-$  collider at  $\sqrt{s_{ee}} = 500$  GeV.

### The two $x$ -ranges studied

LC2:  $0.3 < x < 0.8$

LC1:  $0.1 < x < 0.6$

### The extension of the measurement

- Sys. error = 0.5 OPAL(135 GeV<sup>2</sup>).
- At the ILC the  $Q^2$  range can be extended by about a factor of 40.
- At largest  $Q^2$  this pQCD prediction gets most precise:  
 $\Delta\alpha_s(M_{Z_0}^2)_{\text{theo.}} \rightarrow \mathcal{O}(0.002)$ .

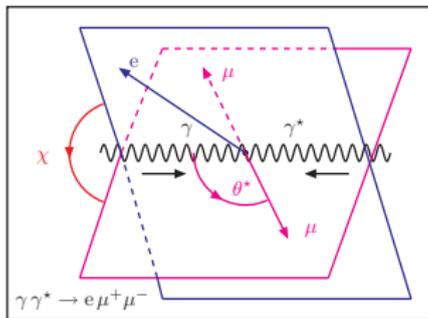
The ILC will help to test this pQCD prediction.

## Conclusions and Outlook

- The structure of the photon has been investigated in great detail by measurements of photon structure functions at  $e^+e^-$  collider.
- The QED structure shows the expected suppression with  $Q^2$ . Also azimuthal correlations and the presence of the interference terms have been observed (both not shown).
- The hadronic structure of the photon is richer than that of the proton due to the presence of both the point-like and the hadron-like components.
- The combined effort of the LEP experiments led to improvements in the description of the hadronic final state by Monte Carlo models.
- The low- $x$  reach of  $F_2^\gamma$  is limited. However the charm contribution to  $F_2^\gamma$  as well as the positive scaling violations of  $F_2^\gamma$  for all  $x$  have been clearly observed.
- A number of new parametrisations of  $F_2^\gamma$  with different theoretical preferences have been obtained in the last years.
- The measurements of  $F_2^\gamma$  should be continued by Babar/Belle and also by the ILC.

**I hope that  $F_2^\gamma$  measurements will be performed at present and future experiments.**

## Azimuthal correlations in muon-pair events



$$-d\sigma \propto 1 - \rho(y) \frac{F_A^\gamma}{F_2^\gamma} \cos \chi + \frac{1}{2} \epsilon(y) \frac{F_B^\gamma}{F_2^\gamma} \cos 2\chi$$

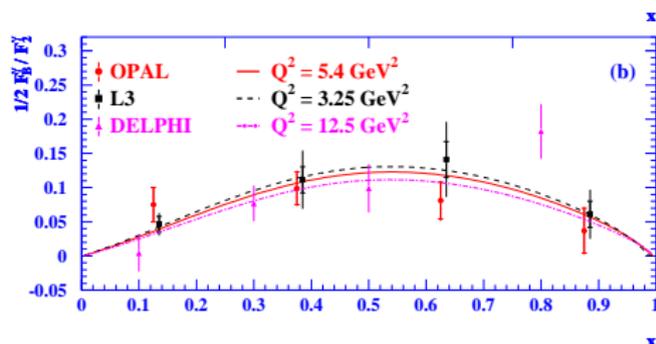
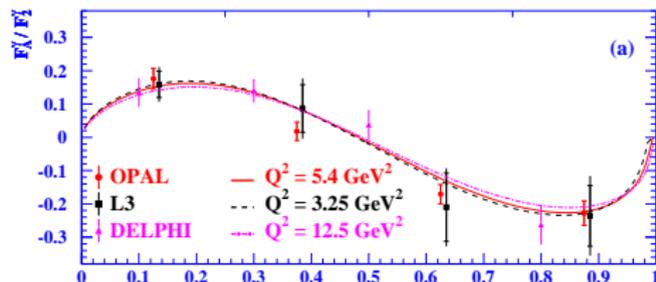
$$\text{with: } \epsilon(y) = \frac{2(1-y)}{1+(1-y)^2} \approx 1, \text{ and}$$

$$\rho(y) = \frac{(2-y)\sqrt{1-y}}{1+(1-y)^2} \approx 1$$

### The probed helicity structure

–  $F_A^\gamma$  transverse-longitudinal interference

–  $F_B^\gamma$  transverse-transverse interference



The  $\chi$  dependence gives access to other structure functions besides  $F_2^\gamma$ .